

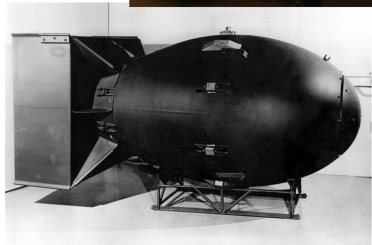
Entrainment of Coupled Phase Oscillators

Jordan Snyder

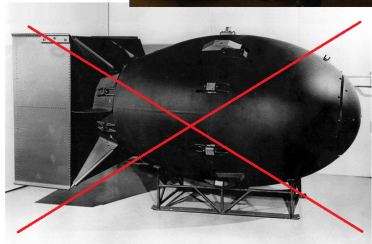


Pizza Seminar, 12 October 2016

Los Alamos



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Outline

- 1 Motivation and Background
 - Collective behavior meets Control
 - Example System
 - Entrainment
 - Coupled Oscillators
- 2 Present Work
 - Forced, Coupled Oscillators
 - Linear Stability Analysis
 - Conclusions
- 3 Future Work
 - Structured Coupling
 - Multiple Timescales



Complexity

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- Global behavior not obvious from local dynamics
- “Magnets”



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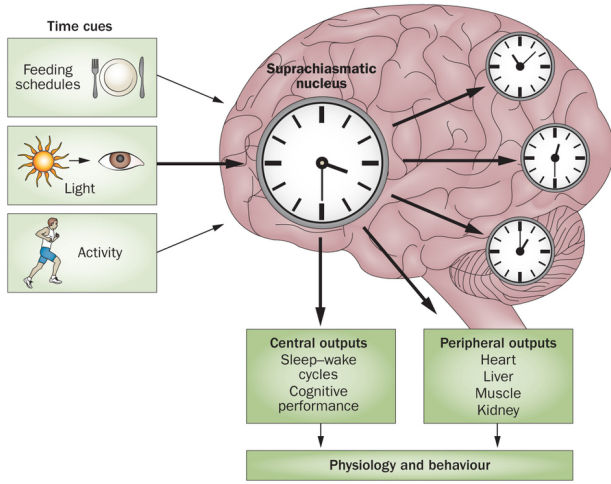
Realistically:

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- As few measurements as possible \implies *open-loop control*

\implies **Goal:** Study a simple system where we can analyze the interplay of open-loop control and collective behavior.

\implies **Q:** Which simple system?

Inspiration : Biological Rhythms



[Videnovic *et al.*, 2014]

Model Class

Desired features:

- Many coupled units with different intrinsic, periodic behavior
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Model System : Population of nonlinear oscillators
subject to *forcing* and *coupling*

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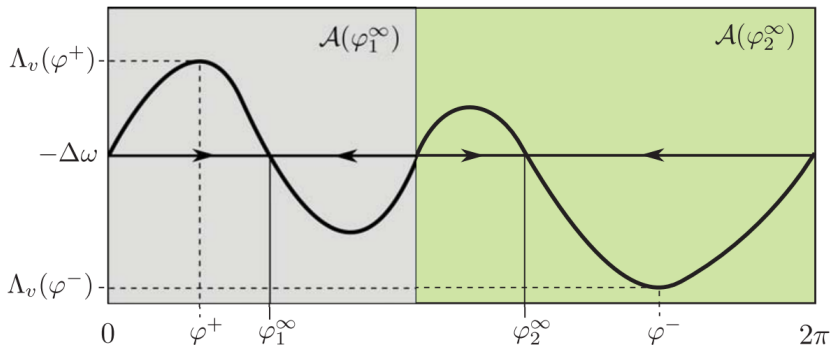
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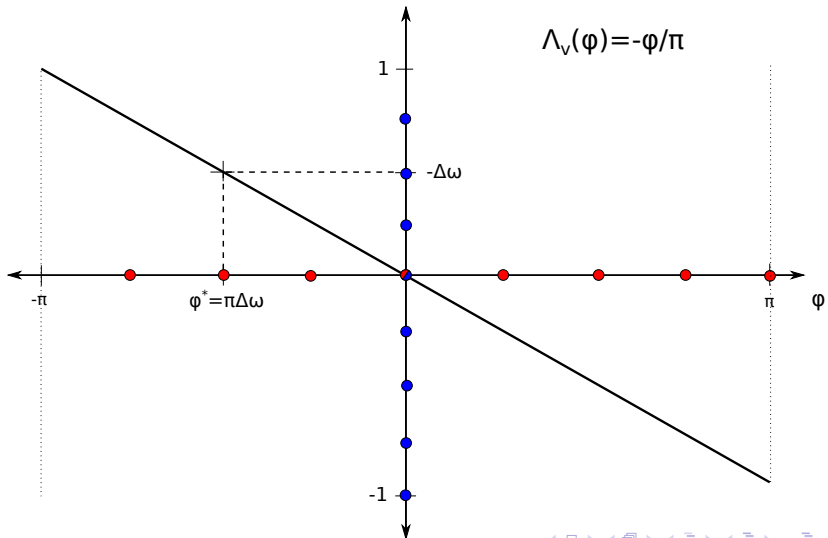
Entrainment

A phase φ^* is a stable fixed point of $\dot{\varphi} = \Delta\omega + \Lambda_v(\varphi)$ if

$$\Delta\omega + \Lambda_v(\varphi^*) = 0 \text{ and } \frac{d\Lambda_v}{d\varphi}(\varphi^*) < 0$$



Example : Entrained Decoherence



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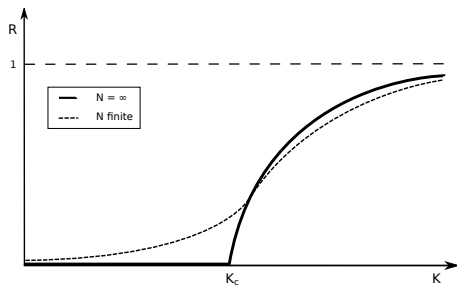
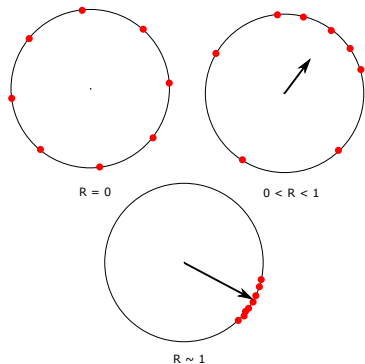
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Trade-off between heterogeneity and coupling at *critical coupling strength* $K_c = \frac{2}{\pi g(0)}$

Phase transition in the Kuramoto model



$K_c =$ point when $R = 0$ state becomes *linearly unstable*

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Entrained Decoherence with Coupling

$$\dot{\varphi}_i = \omega_i + \Lambda_v(\varphi_i) + \frac{K}{N} \sum_{j=1}^N \sin(\varphi_j - \varphi_i)$$

where $\{\omega_i = \frac{2}{N}i - 1\}$ and $\Lambda_v(\varphi) = -\frac{\varphi}{\pi}$ for $\varphi \in [-\pi, \pi)$.

Decoherence - $\{\varphi_i = \pi\omega_i\}$ - is an $R = 0$ fixed point, by symmetry of the coupling term.

Q: For what values of K is this state stable?

We analyze this model for finite N and in the limit $N \rightarrow \infty$

Linear Stability in Finite Dimensions

At the fixed point $\{\varphi_i^* = \pi\omega_i\}$, the Jacobian has matrix elements

$$J_{ij} = \left(\frac{-1}{\pi} - \frac{K}{N} \sum_{k \neq i} \cos(\varphi_k^* - \varphi_i^*) \right) \delta_{ij} + (1 - \delta_{ij}) \frac{K}{N} \cos(\varphi_j^* - \varphi_i^*)$$

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Bound the eigenvalues using the Gershgorin circle theorem:

$$\Re(\lambda) \leq \frac{-1}{\pi} + K \implies K_c \geq \frac{1}{\pi}$$

The limit $N \rightarrow \infty$

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Dynamics are given by a *continuity equation*:

$$\partial_t \rho_\omega + D(v_\omega \rho_\omega) = 0$$

where D is the derivative in the sense of distributions, and $v_\omega = v_\omega(\varphi)$ is the *phase velocity*.

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Linearized dynamics near $\{\rho_\omega = \delta_{\pi\omega}\}$ can be diagonalized exactly:

$$\sigma(L) = \left\{ \frac{-1}{\pi}, \frac{-1}{\pi} + \frac{K}{2} \right\} \implies K_c = \frac{2}{\pi}$$

Conclusions

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Despite phase diversity, external forcing has brought the system closer to order.

Structure \leftrightarrow Function

The interplay between connectivity and dynamics has been studied in many contexts

- Resilience to random breakdown: $1 - p_c = \frac{\langle k \rangle}{\langle k^2 \rangle - \langle k \rangle}$
- Analyzing structure via dynamics: PageRank
- Master Stability Function: MSF : $\sigma(A) \rightarrow \{\text{stable, not stable}\}$

In synchronization, studies have focused on:

- Cluster sizes for lattices in the limit $N \rightarrow \infty$ [Strogatz and Mirollo, 1988]
- Existence & uniqueness of fixed points [Jadbabaie *et al.*, 2004]
- Paths to synchronization on different topologies [Gómez-Gardenes *et al.*, 2007]
- Correlations, e.g. assortativity [Restrepo and Ott, 2014]

Proposed Work

Questions:

- Given coupling topology, what is the change in stability of decoherence upon driving?
- Is it possible to tune the trade-off between driving and coupling by adjusting coupling topology?

Approaches:

- Numerical simulation
- Linear Stability Analysis
- Mean-field Approximation

Subharmonic Entrainment

It is well known that oscillators can behave coherently without attaining the same frequency. Examples include:

- 4-5 day rodent estrous cycle locking to night-day cycle [Winfree, 2001]
- 2:1 phase locking of body temperature and sleep-wake cycle when deprived of external time cues [Aschoff and Wever, 1981]
- Subharmonic locking of a sensory neuron to periodic inhibitory input [Perkel *et al.*, 1964]

Subharmonic entrainment of a single oscillator is well described mathematically (in, e.g., [Zlotnik and Li, 2014]). However, coupling of different subharmonically forced oscillators is not well understood.

Proposed Work

Questions:

- Can common driving selectively help or hinder mutual $N: M$ entrainment?
- Can coupling across timescales improve coherence within timescales?

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Pie-in-the-sky scientific question... Is the seven-day week adaptive?

References I

Jürgen Aschoff and Rütger Wever. The circadian system of man. In *Biological rhythms*, pages 311–331. Springer, 1981.

Florian Dörfler and Francesco Bullo. Synchronization in complex networks of phase oscillators: A survey. *Automatica*, 50(6):1539–1564, 2014.

Jesús Gómez-Gardenes, Yamir Moreno, and Alex Arenas. Paths to synchronization on complex networks. *Physical review letters*, 98(3):034101, 2007.

Ali Jadbabaie, Nader Motee, and Mauricio Barahona. On the stability of the kuramoto model of coupled nonlinear oscillators. In *American Control Conference, 2004. Proceedings of the 2004*, volume 5, pages 4296–4301. IEEE, 2004.

Hiroya Nakao. Personal Website.

References II

- Donald H Perkel, Joseph H Schulman, Theodore H Bullock, George P Moore, and Jose P Segundo. Pacemaker neurons: effects of regularly spaced synaptic input. *Science*, 145(3627):61–63, 1964.
- Juan G Restrepo and Edward Ott. Mean-field theory of assortative networks of phase oscillators. *EPL (Europhysics Letters)*, 107(6):60006, 2014.
- Steven H Strogatz and Renato E Mirollo. Collective synchronisation in lattices of nonlinear oscillators with randomness. *Journal of Physics A: Mathematical and General*, 21(13):L699, 1988.
- Aleksandar Videnovic, Alpar S Lazar, Roger A Barker, and Sebastiaan Overeem. 'the clocks that time us'[mdash] circadian rhythms in neurodegenerative disorders. *Nature Reviews Neurology*, 10(12):683–693, 2014.
- Arthur T Winfree. *The geometry of biological time*, volume 12. Springer Science & Business Media, 2001.

References III

Anatoly Zlotnik and Jr-Shin Li. Optimal subharmonic entrainment of weakly forced nonlinear oscillators. *SIAM Journal on Applied Dynamical Systems*, 13(4):1654–1693, 2014.

Anatoly Zlotnik, Raphael Nagao, István Z Kiss, and Jr-Shin Li. Phase-selective entrainment of nonlinear oscillator ensembles. *Nature communications*, 7, 2016.

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Start with a forced, nonlinear oscillator with stable limit cycle γ :

$$\dot{x} = f(x; u), \quad \dot{\gamma} = f(\gamma; 0), \quad \gamma(t+T) = \gamma(t)$$

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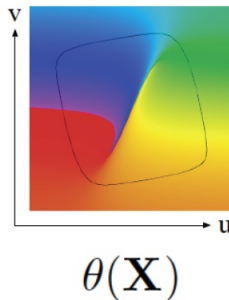
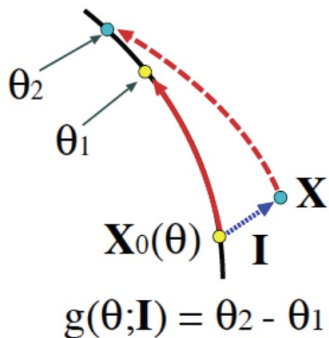
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Forcing and Coupling

Assuming that oscillators *respond* to forces according to $Z(\psi)$, *exert* forces according to $P(\psi)$, and are all subject to the same force $u = v(\Omega t)$,

$$\dot{\psi}_i = \omega_i + Z(\psi_i)v(\Omega t) + \frac{K}{N} \sum_{j=1}^N Z(\psi_i)P(\psi_j)$$

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Changing coordinates, $\psi = \phi + \Omega t$, and averaging over $\Omega t \in [0, 2\pi)$:

$$\dot{\phi}_i = \Delta\omega_i + \Lambda_v(\phi_i) + \frac{K}{N} \sum_{j=1}^N G(\phi_j - \phi_i)$$

“Separate but Comparable”

Separation of timescales underlies much analysis:

- Center manifold reduction
- Model fast degrees of freedom as stationary noise: replace high-dimensional ODE by low-dimensional SDE